
Physics Insight from Plasma Shaping of TCV Tokamak Plasmas - focus on Electron Heat Transport

(TCV = Tokamak à Configuration Variable)

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OUTLINE

1. Motivations for plasma shaping
2. TCV tokamak facility
3. Confinement and electron heat transport vs. shape
 - 3a. Confinement shape enhancement factor (OH, high collisionality)
 - 3b. Electron heat transport (ECH, low collisionality)
4. Conclusions on heat transport
5. Outlook on innovative divertors / shapes

1. MOTIVATIONS of SHAPE STUDIES

WHY STUDY also SHAPES different from ITER?

- Test and validation of MHD and transport theory
- New: negative triangularity *improves confinement* (in L-mode!)
- Confinement of τ_E , n_e , β , fast ions... scales with plasma current I_p and $I_{p \max}$ can be increased by plasma cross-section shaping *without increasing the magnetic field*
- Many other parameters depend on plasma shaping
- Reciprocally, active plasma shaping offers a way to control plasma parameters
- What plasma shape, what divertor for a device beyond ITER?
to maximize performance, and reduce heat loads

Shaping variables

- elongation κ
- triangularity δ , *including negative (D- and inverse-D-shaped)*
- squareness
- limited / various diverted, innovative geometry
- single plasmas / doublets plasmas ...
- beyond TCV: aspect ratio R/a

TCV: $R = 0.88 \text{ m}$, $a = 0.25 \text{ m}$, $R/a \sim 3.5$
 $B < 1.54T$,
 $I_p \leq 1\text{MA}$
elongation $0.9 < \kappa < 2.8$
triangularity - $0.7 < \delta < 1$
(reached parameters)

Many parameters, effects, are influenced by plasma shape

Direct influence of shape ...

- **MHD stability** sawteeth, modes, NTM, disruptions, TAEs, ELMs, ...
- **Transport** electron heat, momentum (particle...)
- **Confinement** pressure limits, edge transport barrier, performance

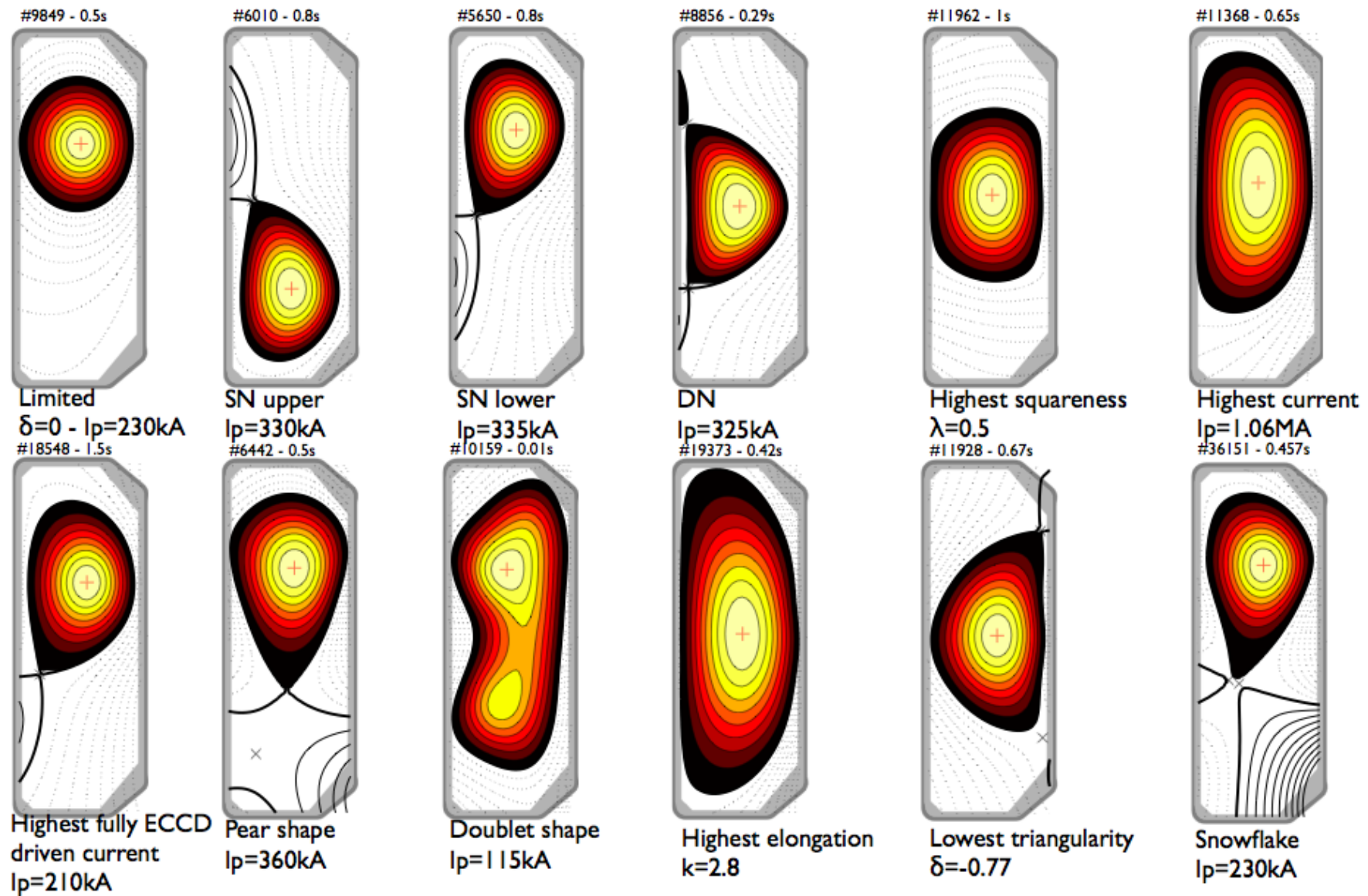
Indirect influence ...

- **ELMs(shape)** [can destroy ITBs (e.g. JET), etc...]
- **Sawteeth(shape)** [can trigger NTMs]

*Dependence on the plasma **confinement regime** ...*

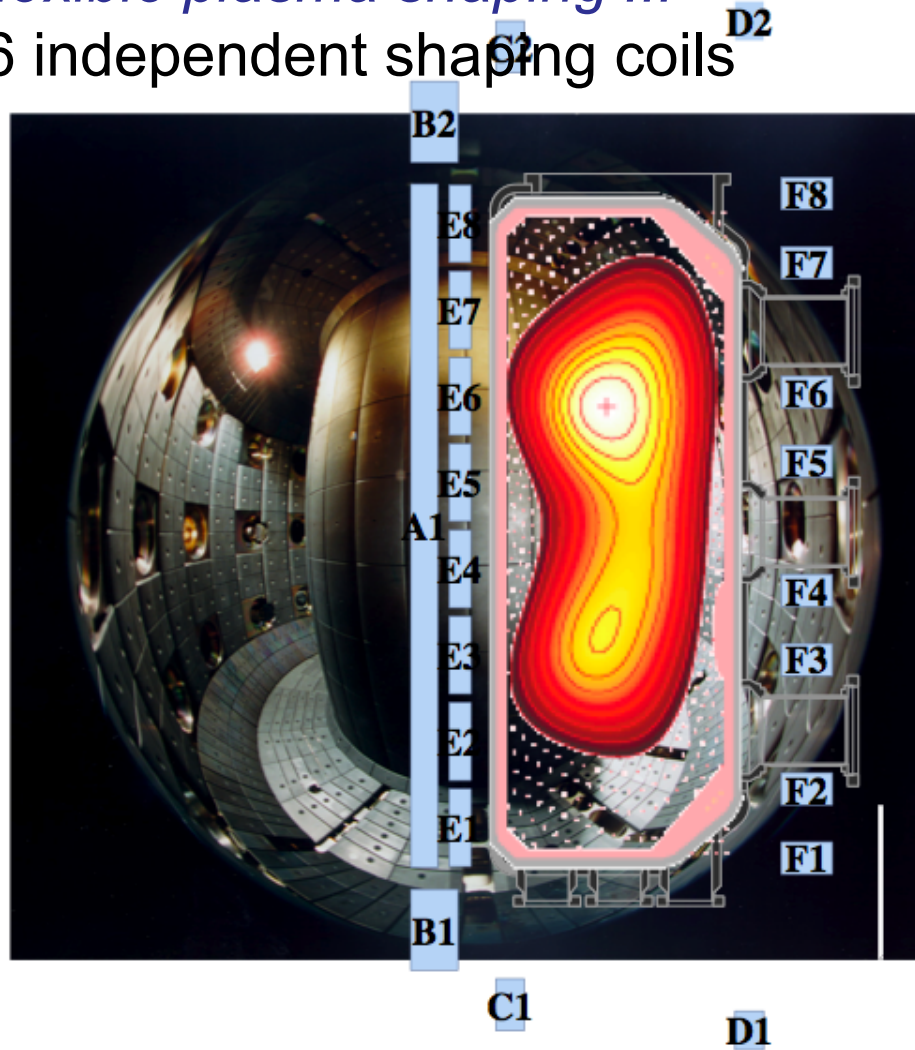
- $\tau_E(\delta)$ increases towards **negative** δ in L-mode (through core transport)
increases towards **positive** δ in H-mode (through pedestal height)

2. TCV FACILITY and SHAPING ACHIEVEMENTS

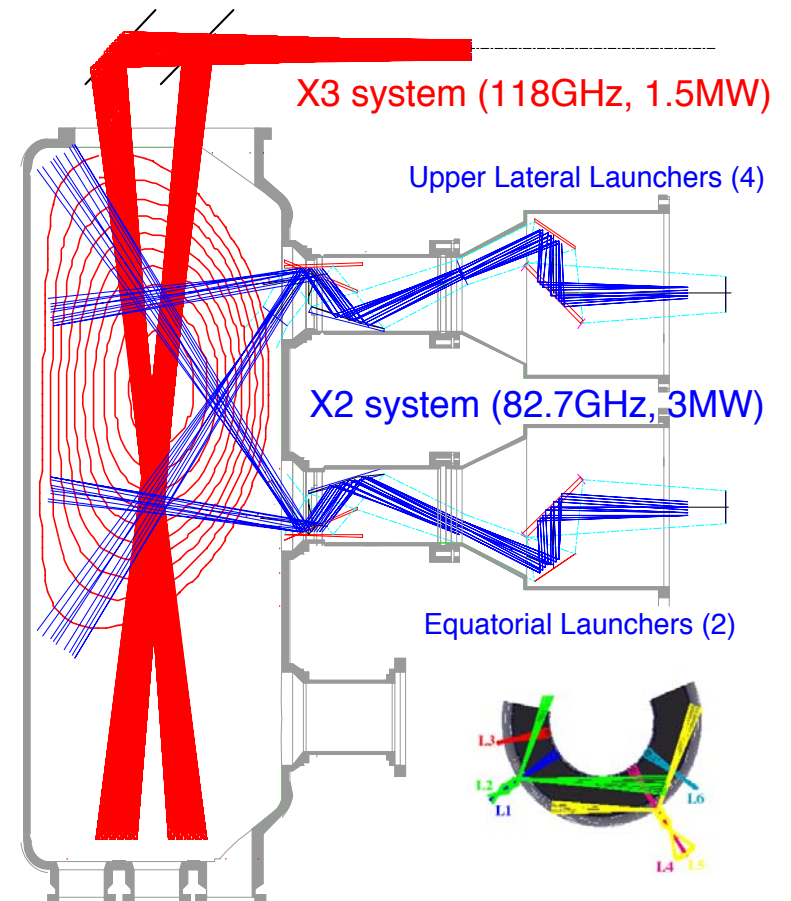


TCV FACILITY

Flexible plasma shaping ...
16 independent shaping coils



... matched by a flexible heating system, entirely based on ECRH



Total: 4.5 MW at 2nd and 3rd harmonic
Cut-off densities: 4.2 and $11.5 \cdot 10^{19} \text{ m}^{-3}$

3a. CONFINEMENT vs. SHAPE / Ohmic, high collisionality

densities ($\nu_{\text{eff}} \sim 2.5-10$)

- τ_{Ee} increases strongly with κ
- Mild decrease with δ , in $\delta > 0$ -range
- Interpretation in terms of a confinement **Shape Enhancement Factor (SEF)**:

$$Q_\alpha = -n_\alpha \chi_\alpha \left\langle |\vec{\nabla} \rho|^2 \right\rangle \frac{\partial T_\alpha}{\partial \rho}$$

heat diffusivity:
depends on collisions,
shape, etc...

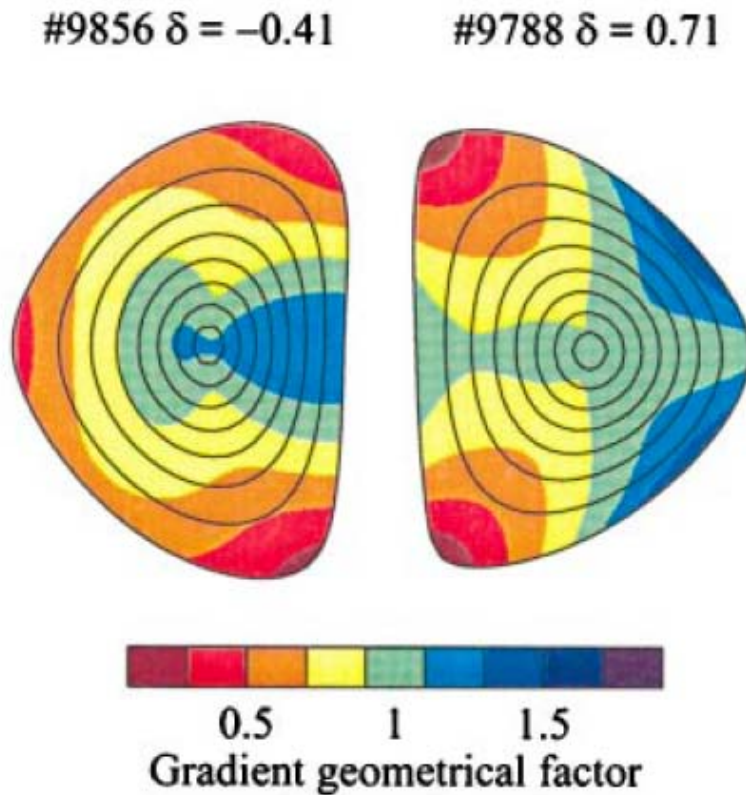
gradient geometrical factor,
expresses flux surfaces sep.,
metric term

Moret PRL97, Weisen NF97

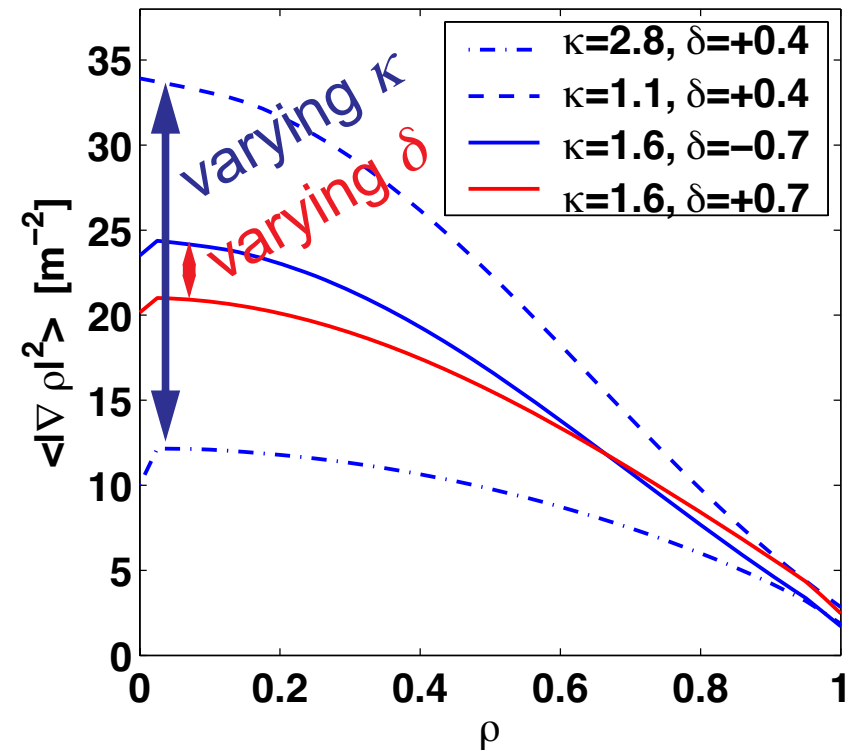
Gradient geometrical factor

gradient geometrical factor

local



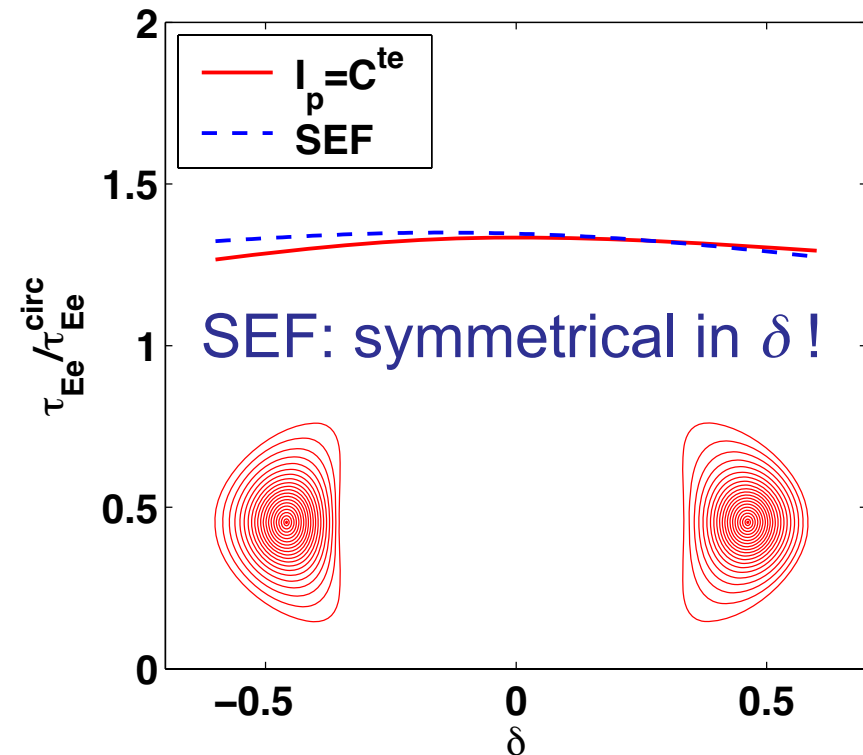
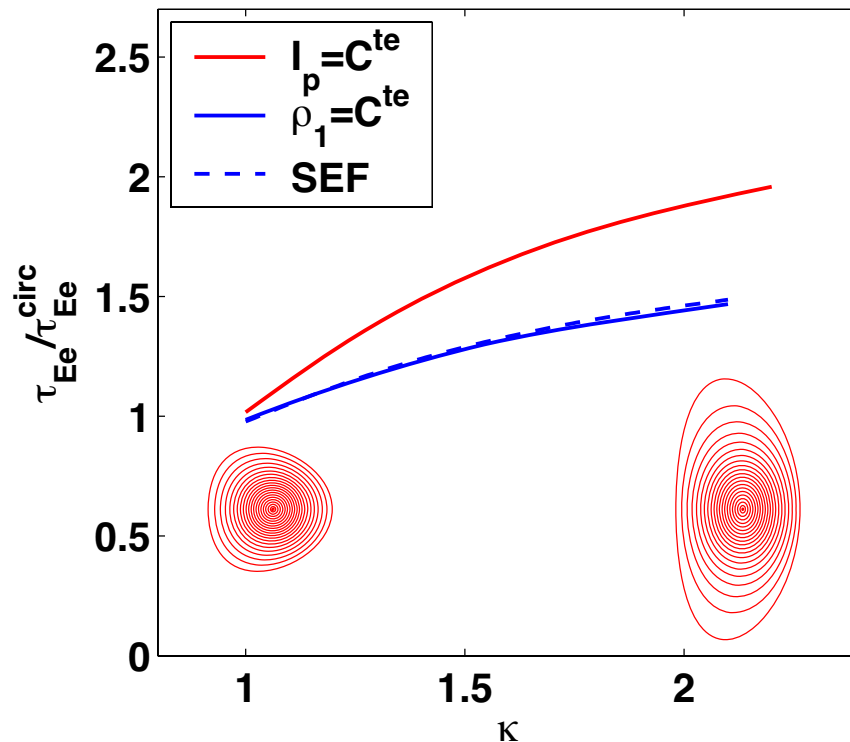
flux surface averaged



Moret PRL97

SEF: when confinement only due to flux surface separation ($\chi_e = \text{cst}$)

τ_{Ee} shaped plasma / τ_{Ee} circular plasma with same χ_e (ASTRA calculations)



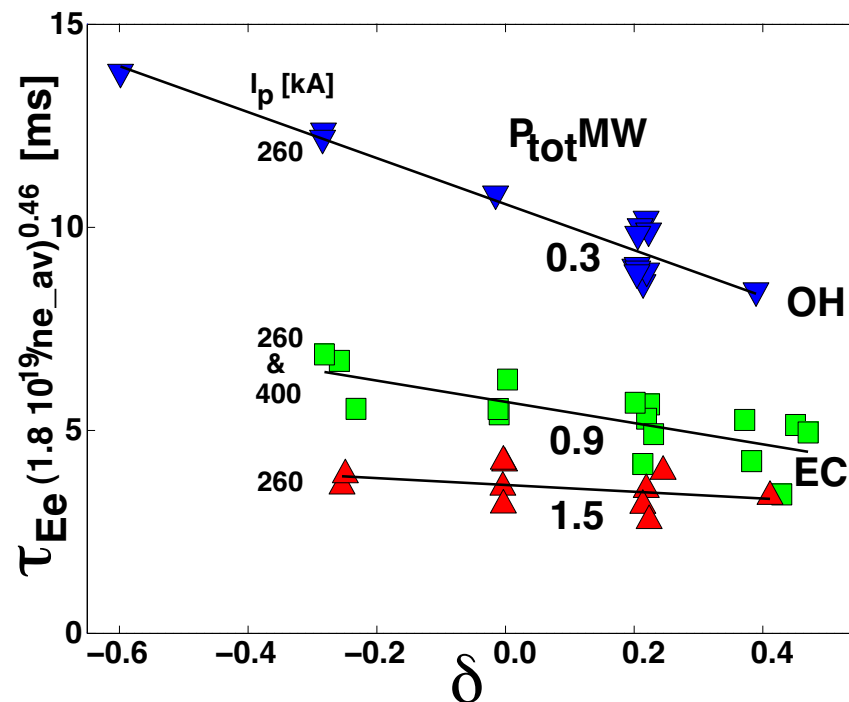
Important to keep sawtooth inv. radius $\rho_1 \sim \text{const}$ (self-similar profiles)

SEF found adequate to account for τ_{Ee} variations with shape in OH, at medium densities

Camenen thesis06

ECH confinement at low densities ($\nu_{\text{eff}} \sim 0.2-1$)

Central ECH, large δ -range, positive and negative triangularity

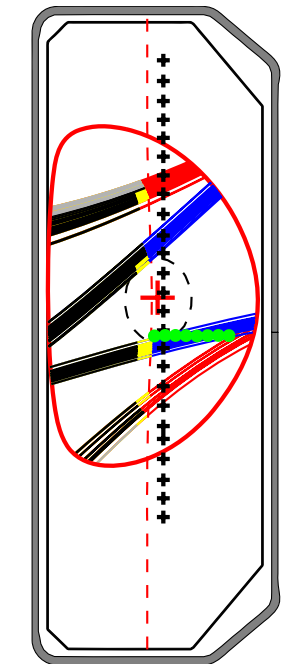


- Strong $\tau_{Ee}(\delta)$ dependence, improvement towards $\delta < 0$
- Asymmetry in triangularity not explained by SEF alone:
—> *does also χ_e vary with shape?*

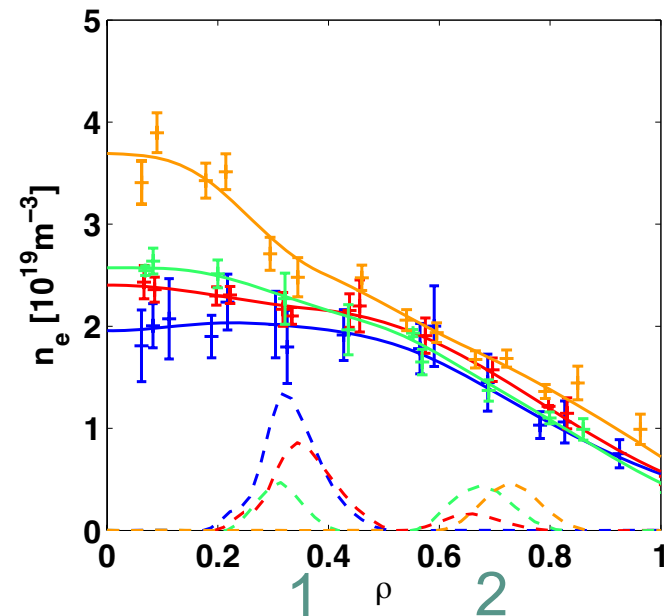
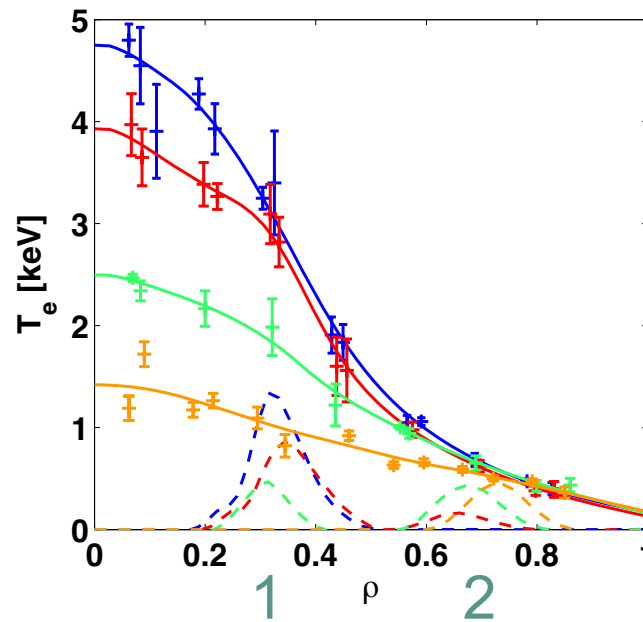
Coda 98, Pochelon NF99, EPS99, Weisen NF98

3b. ELECTRON HEAT TRANSPORT vs shape / EC low n_e ($\nu_{\text{eff}} \sim 0.2-1$)

T_e -variation, $\text{grad} T_e$ -variation experiments



+ diff. Thom.
• CXRS

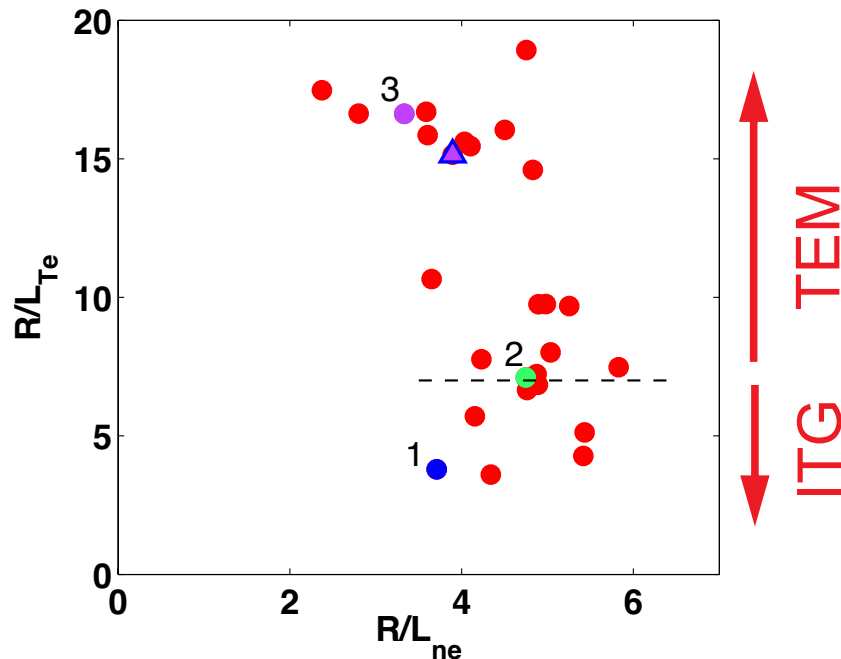


- 2 deposition locations 1 and 2
- Varying $P_{\text{tot}} = P_1 + P_2$
and P_1 / P_2

Camenen PPCF05

Microinstability types in EC plasmas

Type of micro-instabilities (GLF, LORB)



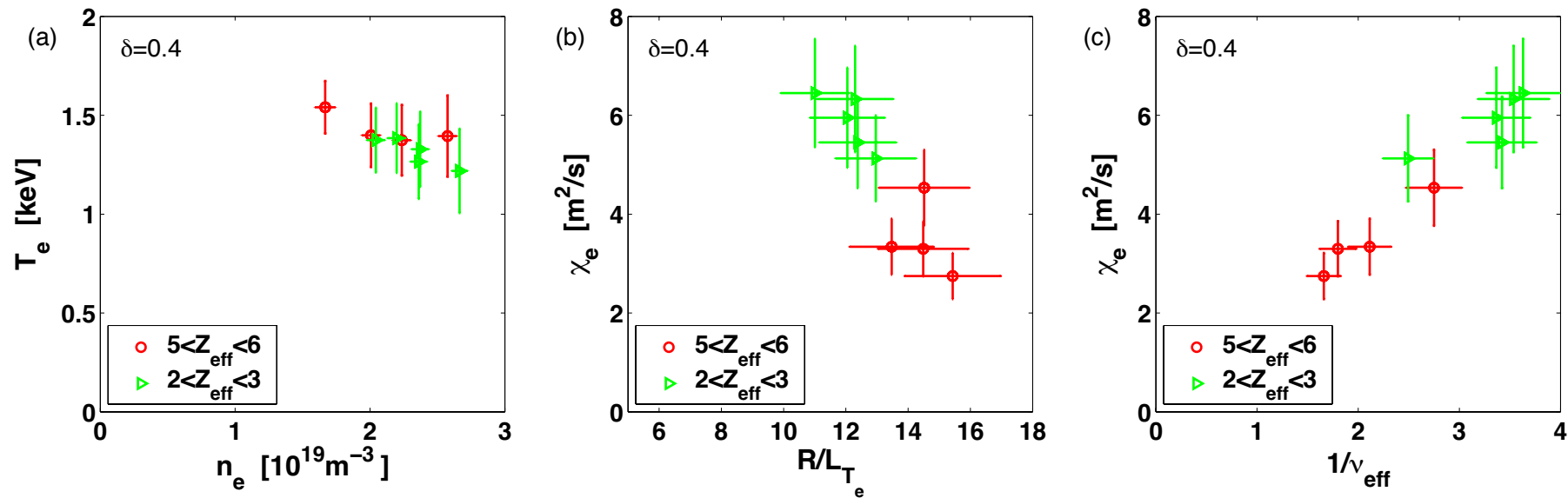
EC plasmas:
High $R/L_{Te} > 7$, in high T_e/T_i ,
low collisionality:

TEM dominated regime

(no ETG in $0.2 < \rho < 0.7$ range,
due to high Z_{eff} & high T_e/T_i)

Camenen PPCF05, Thesis06

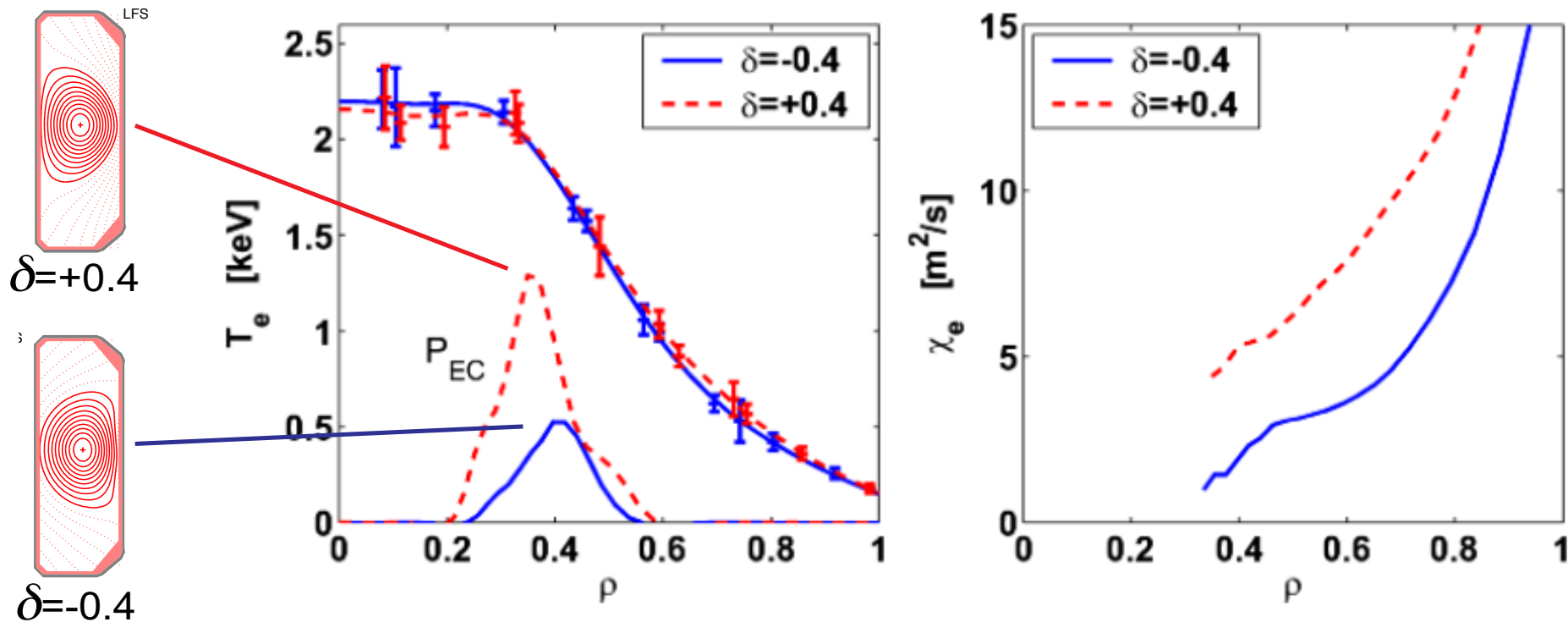
Collisionality effect demonstrated



- The effect of T_e , n_e , Z_{eff} , when combined show a clear dependence of χ_e on collisionality v_{eff}
- $v_{\text{eff}} = 0.1 R n_e Z_{\text{eff}} / T_e^2$
 $= v_{ei} / \omega_{De}$ De: curvature drift
- Diffusivity χ_e reduces with increasing collisionality

Camenen NF07

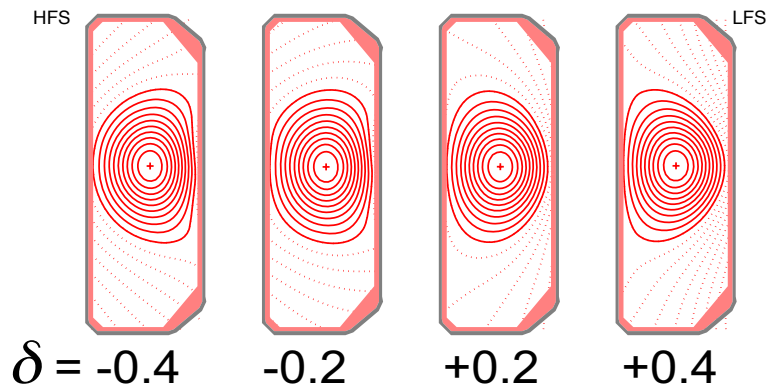
How to decrease transport by a factor 2 in L-mode ?



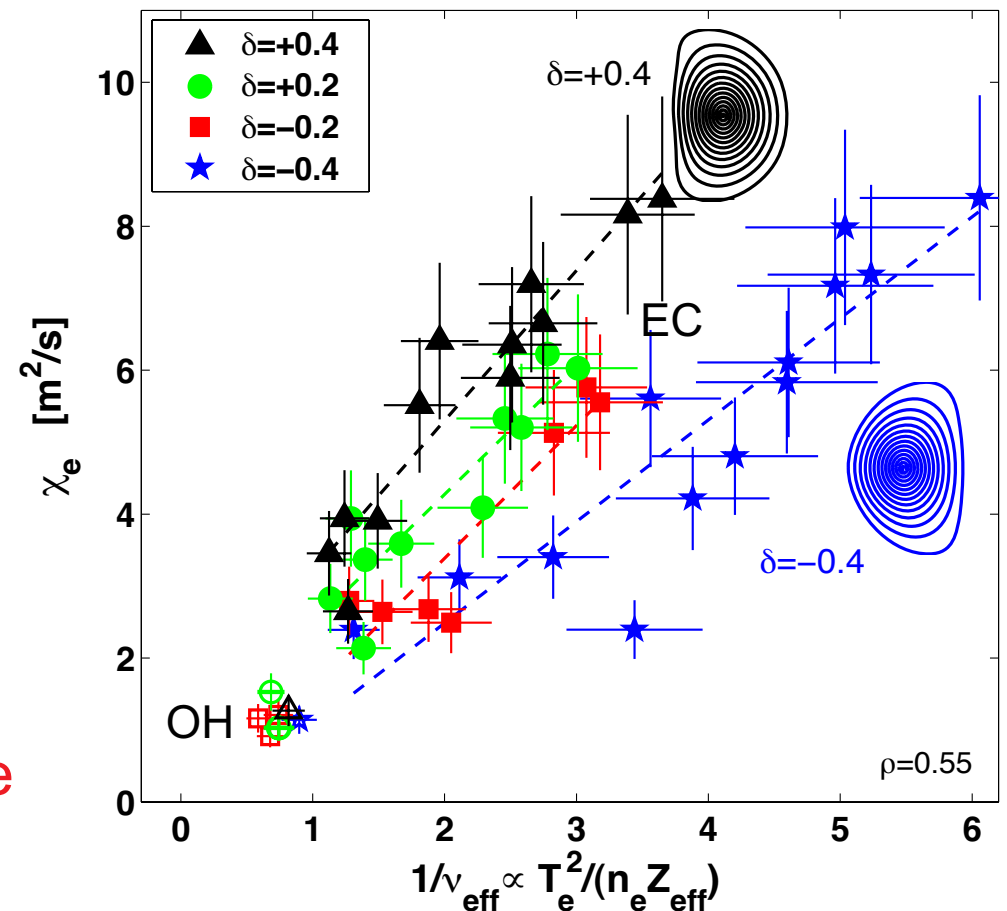
- 1) change triangularity from positive to negative :
 $\delta = +0.4 \rightarrow -0.4$
- 2) adapt additional power to keep the same T_e , n_e and q -profiles
 \rightarrow Result: heat transport is **halved** at mid-radius,
an effect of **negative triangularity**

Camenen NF07

Shape and collisionality effects separated



- Collisionality lowers χ_e
- At same ν_{eff} , negative δ reduces χ_e
- At high ν_{eff} , vanishing δ -effect
- Taking account of collisionality ν_{eff} explains the different δ -dependence of OH and EC data

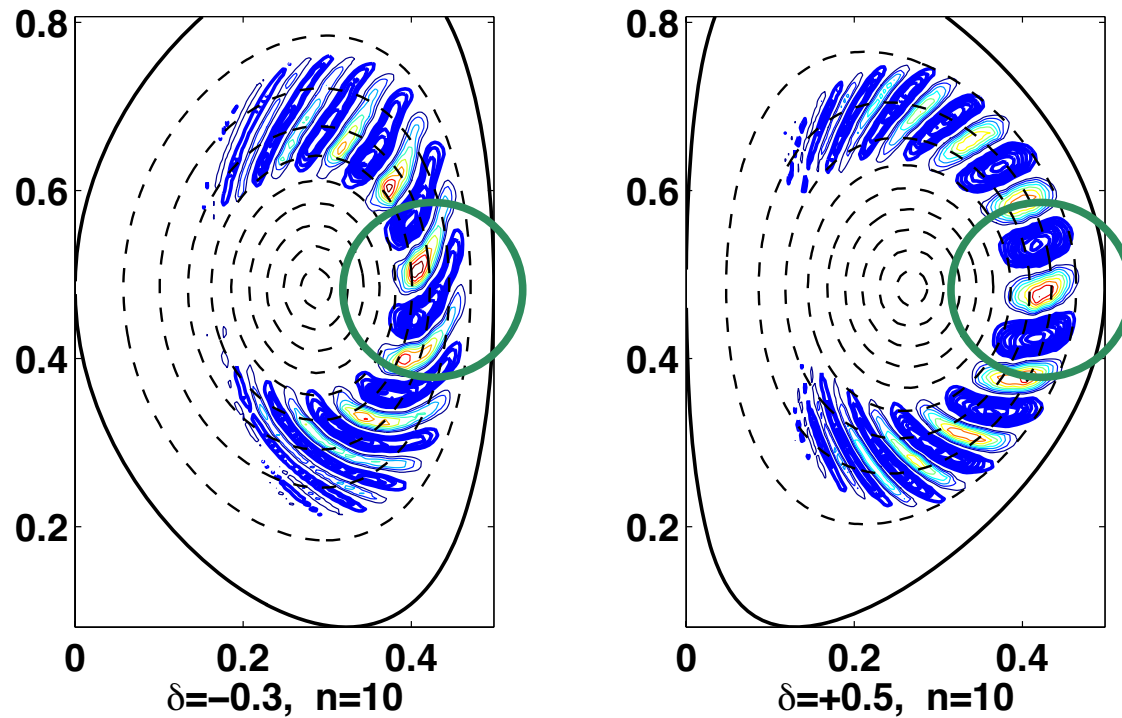


Camenen NF07

Global TEM micro-instabilities simulations

LORB simulations (gyrokinetic, linear, global, no collisions)

Camenen PPCF05



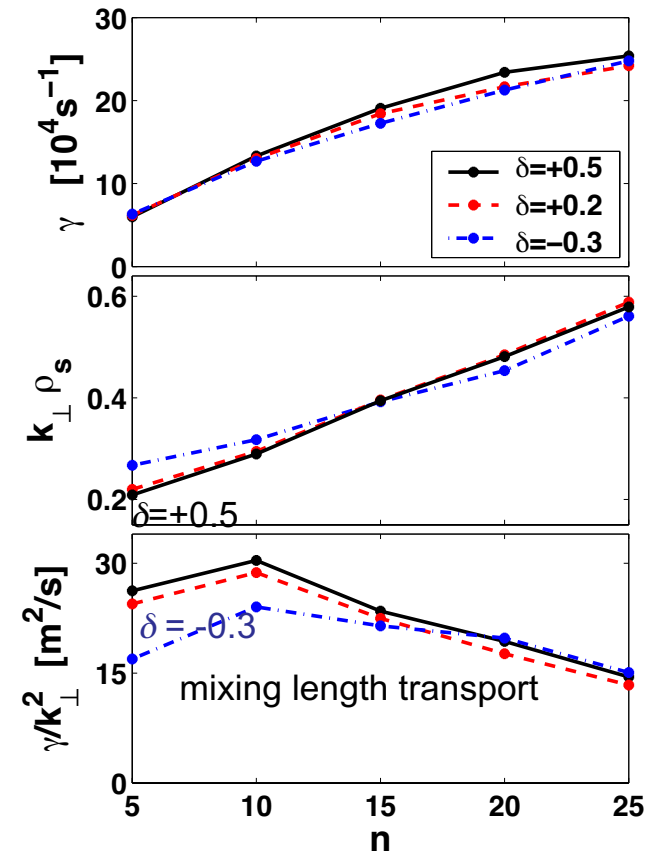
electrostatic potential fluctuations

Negative $\delta < 0$:

-> stronger tilt of eddies at LFS equator!

-> higher k_{\perp}

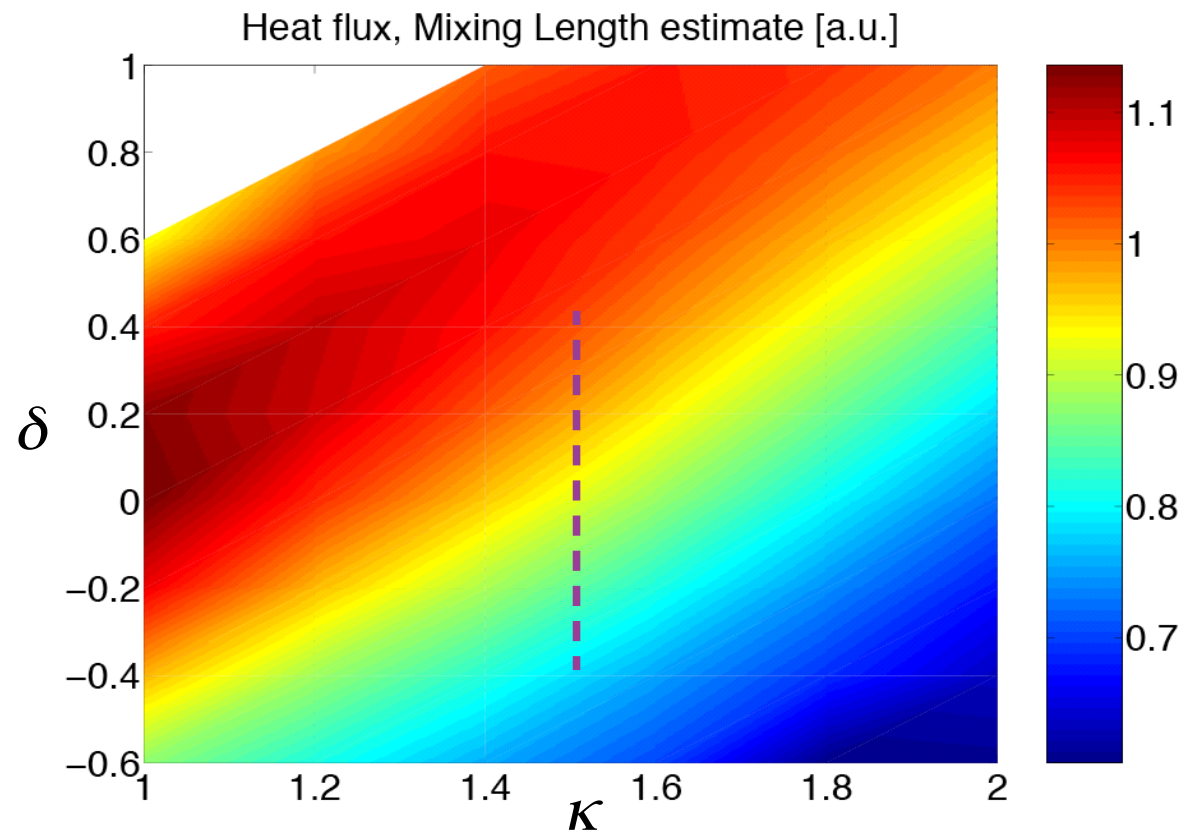
-> lower mixing-length transport



GS2 linear, local, gyrokinetic calculations

χ_e mixing length as a function of (κ, δ)

GS2 (gyrokinetic, linear, local)



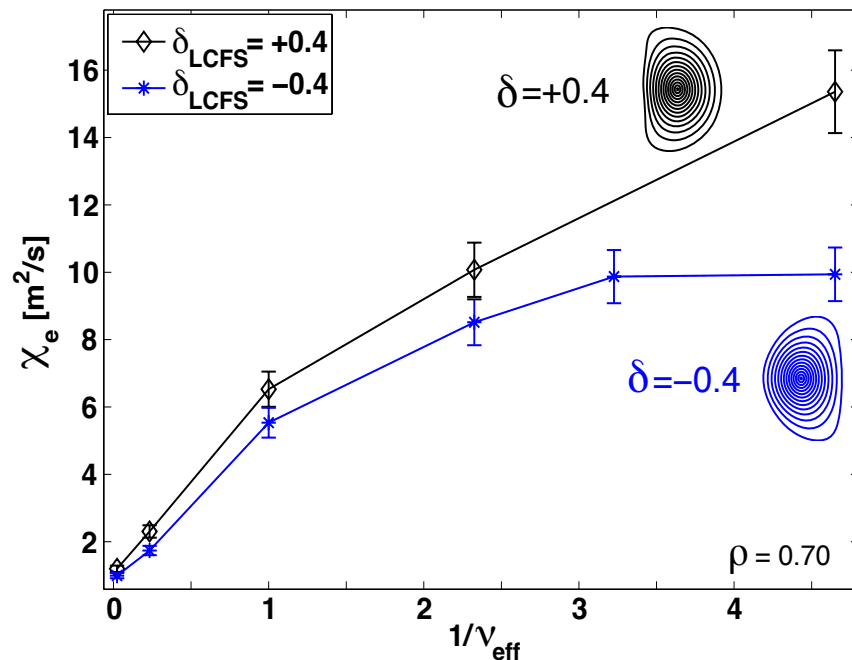
As in above δ -scan expts,
 χ_{e_GS2} reduces monotonically
towards $\delta < 0$

Marinoni PPCF 09

GS2: local, non-linear gyrokinetic calculations, with collisions

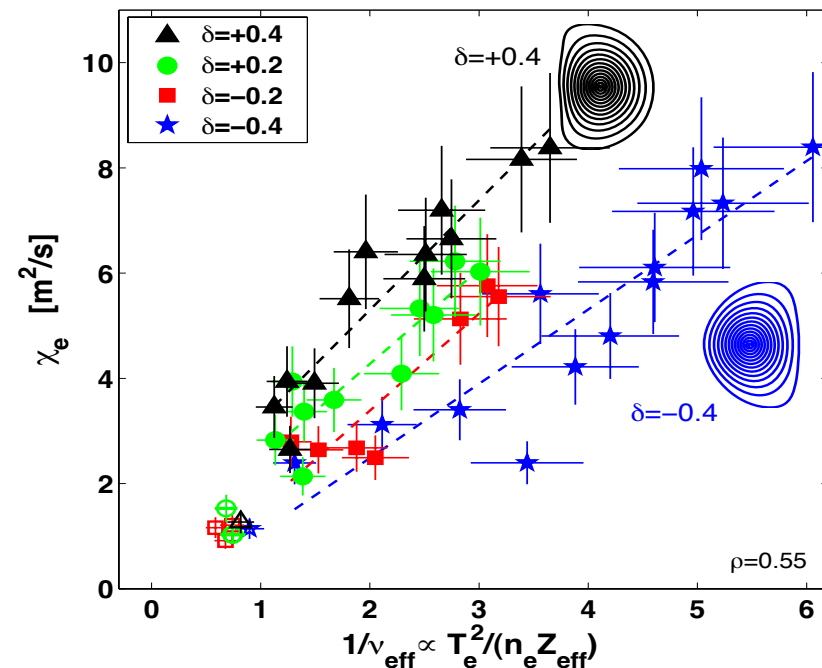
effect of shape and collisionality

GS2



Marinoni PPCF 09

TCV

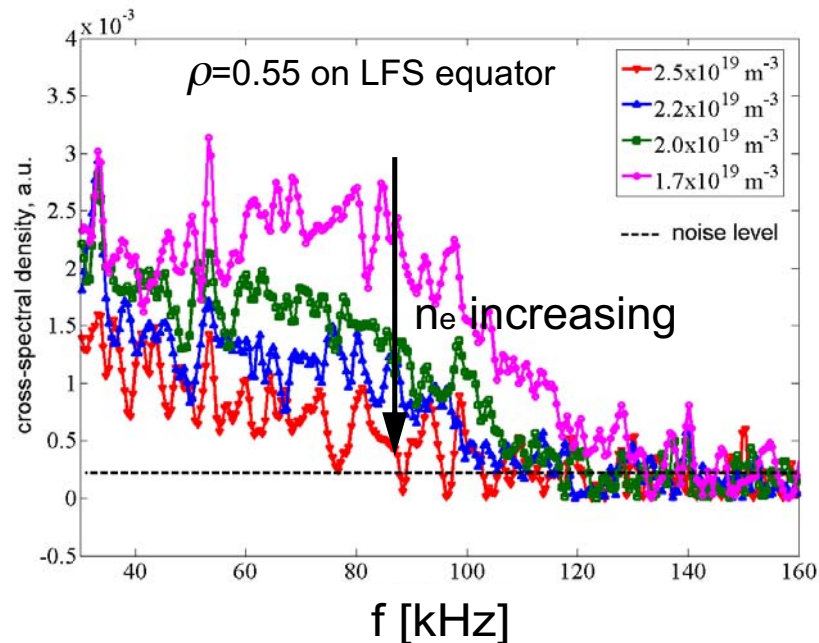


Transport simulations reflect experimental χ_e in TEM regime:

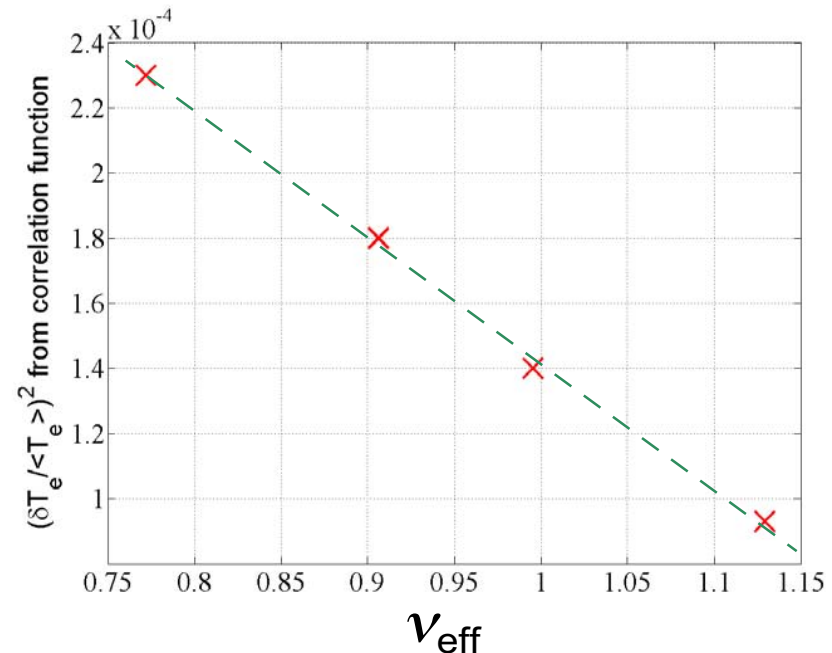
- decrease of χ_e towards high ν_{eff} and negative δ
- triangularity effect on χ_e smaller at high ν_{eff}
- but disagree for the radial dependence: possibly a global effect.

T_e -fluctuations decrease with collisionality (corr-ECE diag.)

fluctuation spectra decrease with density
Ohmic, $q \sim 10$, $\kappa \sim 1.4$, $\delta \sim 0.3$



fluctuation amplitude $(\delta T_e / \langle T_e \rangle)^2$
averaged over $<30-130 \text{ kHz}>$



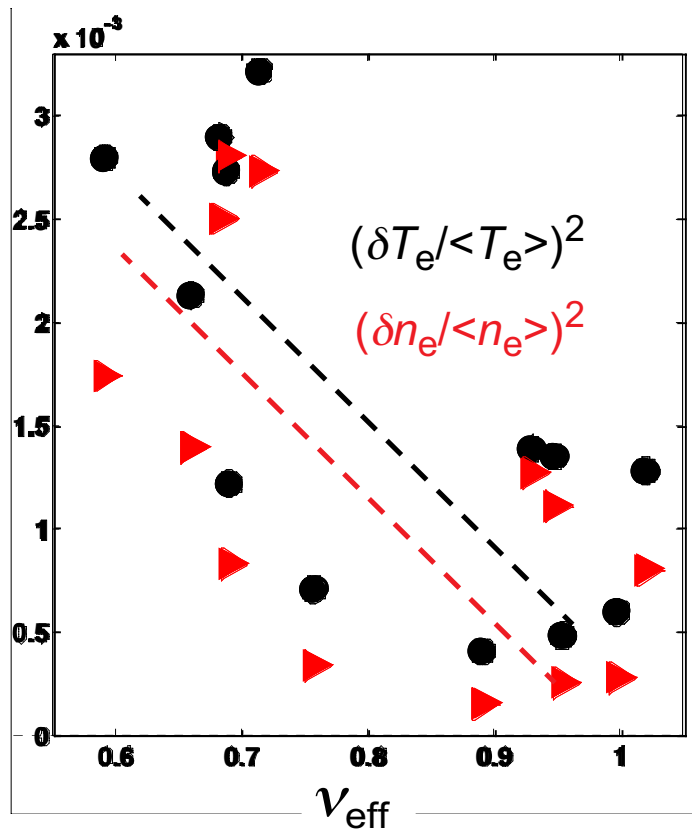
T_e -fluctuations amplitude decrease with collisionality ν_{eff}

Udintsev & Fable US-TTF09

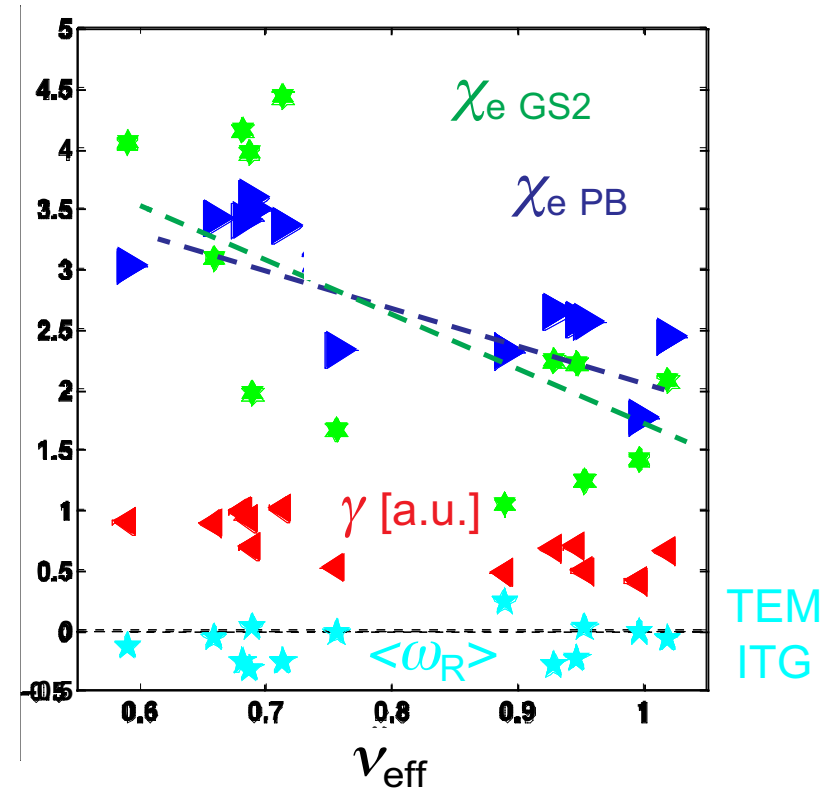
de Meijere: talk 15h (534, in place of 504)
More on turbulence measurements in TCV

Fluctuation amplitude decrease with ν_{eff} , as predicted by GS2

T_e - & n_e -fluctuations from GS2 decrease with ν_{eff}
(as in experiment)



χ_{e_GS2} decreases with ν_{eff} ,
as χ_{e_PB} from experiment

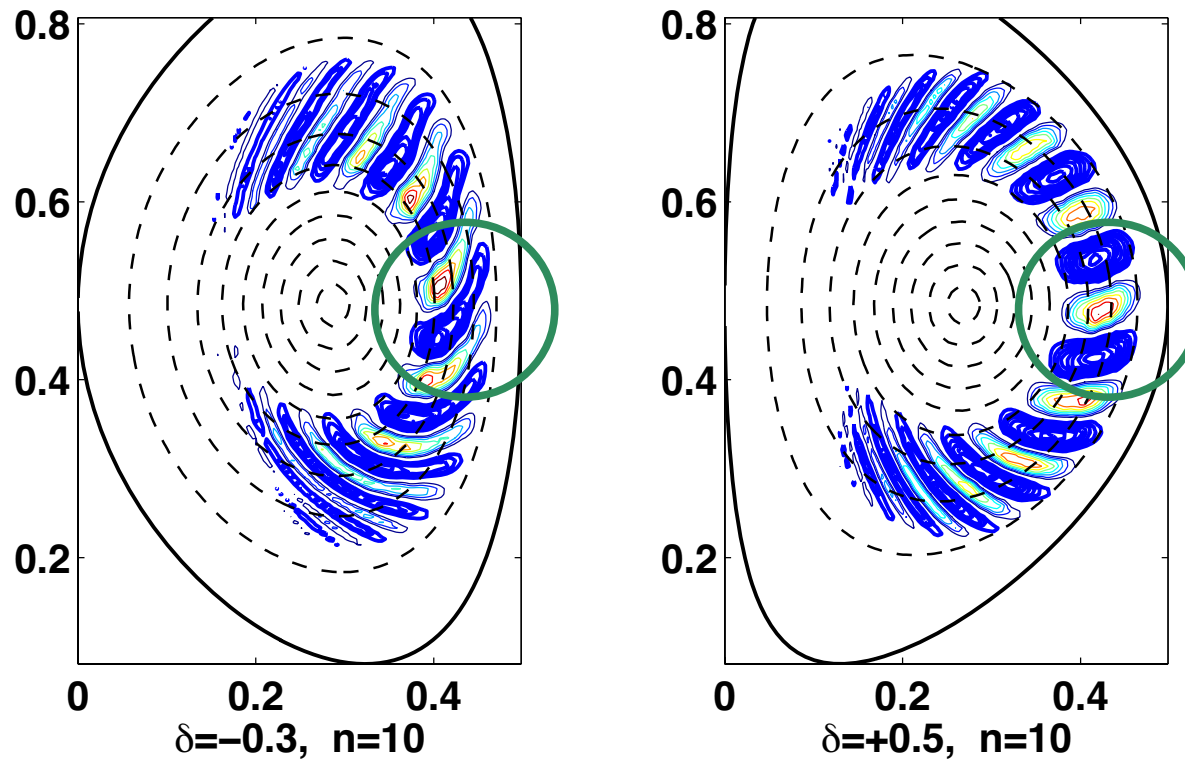


consistent with TEM ampl. reduction with collisions (e^- collisional detrapping)

Udintsev, Fable US-TTF09

Evaluating turbulence correlation length with «corr-ECE diag»

Change of eddy geometry with triangularity
to be measured by correlation-ECE measurements



*How close/different are the experimental results
to the linear (or non-linear) global gyrokinetic calculations?*

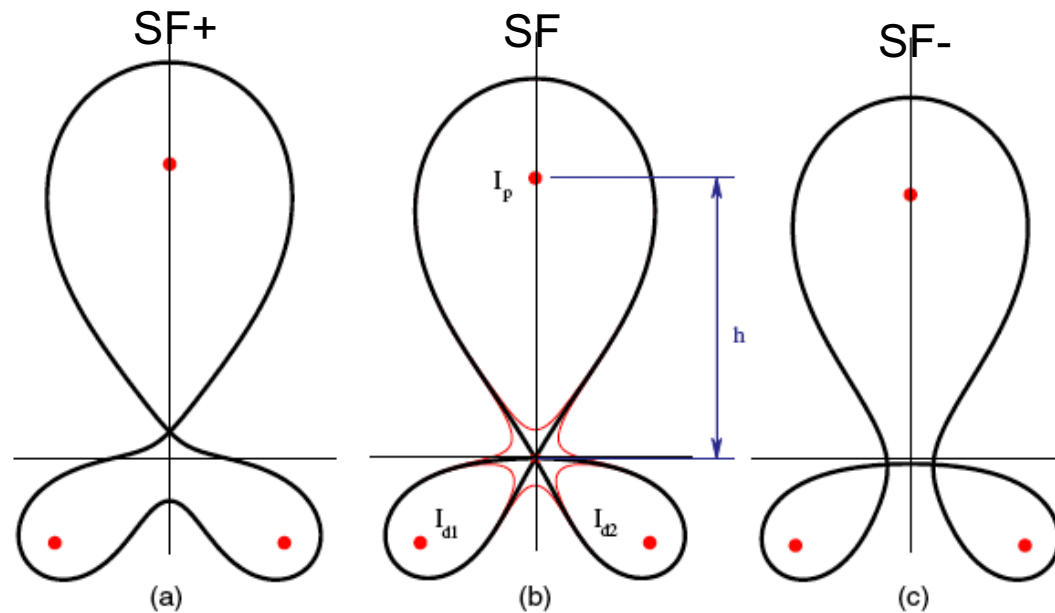
4. CONCLUSIONS on heat transport / shape

- **Role of metric** (flux surface separation)
SEF, evidenced at high ν_{eff} (where χ_e weakly dependent on shape)
Confinement improves with elongation
- **Role of collisionality** (χ_e)
Effect of collisionality evidenced varying triangularity (where $\text{SEF} \sim \text{const.}$)
At low ν_{eff} , τ_e and χ_e are not found symmetric in δ ,
Confinement improving towards negative δ (in TEM dom. transp.)
- **TEM dominated transport**
 - χ_e decreases with ν_{eff} (electron de-trapping)
 - $\langle \delta T_e / T_e \rangle^2$ decreases with ν_{eff}
 - role of triangularity in (de)tuning resonance between ω_D and ω_{TEM}
 - global GK calculations show larger k_{\perp} at $\delta < 0$, thus smaller $\chi_{e_ml} \sim \gamma / k_{\perp}^2$
 - an effect of magn. curvature & Shafranov shift
 - to be identified at the turb. level (e.g. corr-ECE)

5. OUTLOOK on INNOVATIVE SHAPES: 5.1 Snowflake divertor

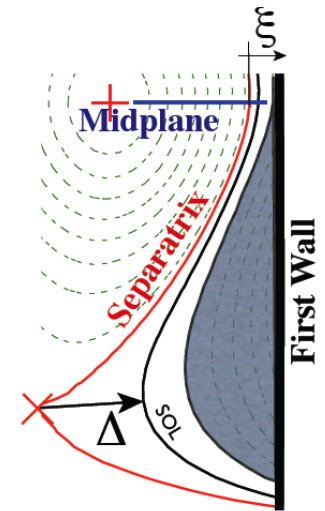
On the search of reduction of divertor power heat loads:

- standard quadrupole null \rightarrow hexapole null (= SF divertor)



Ryutov PP 07 08

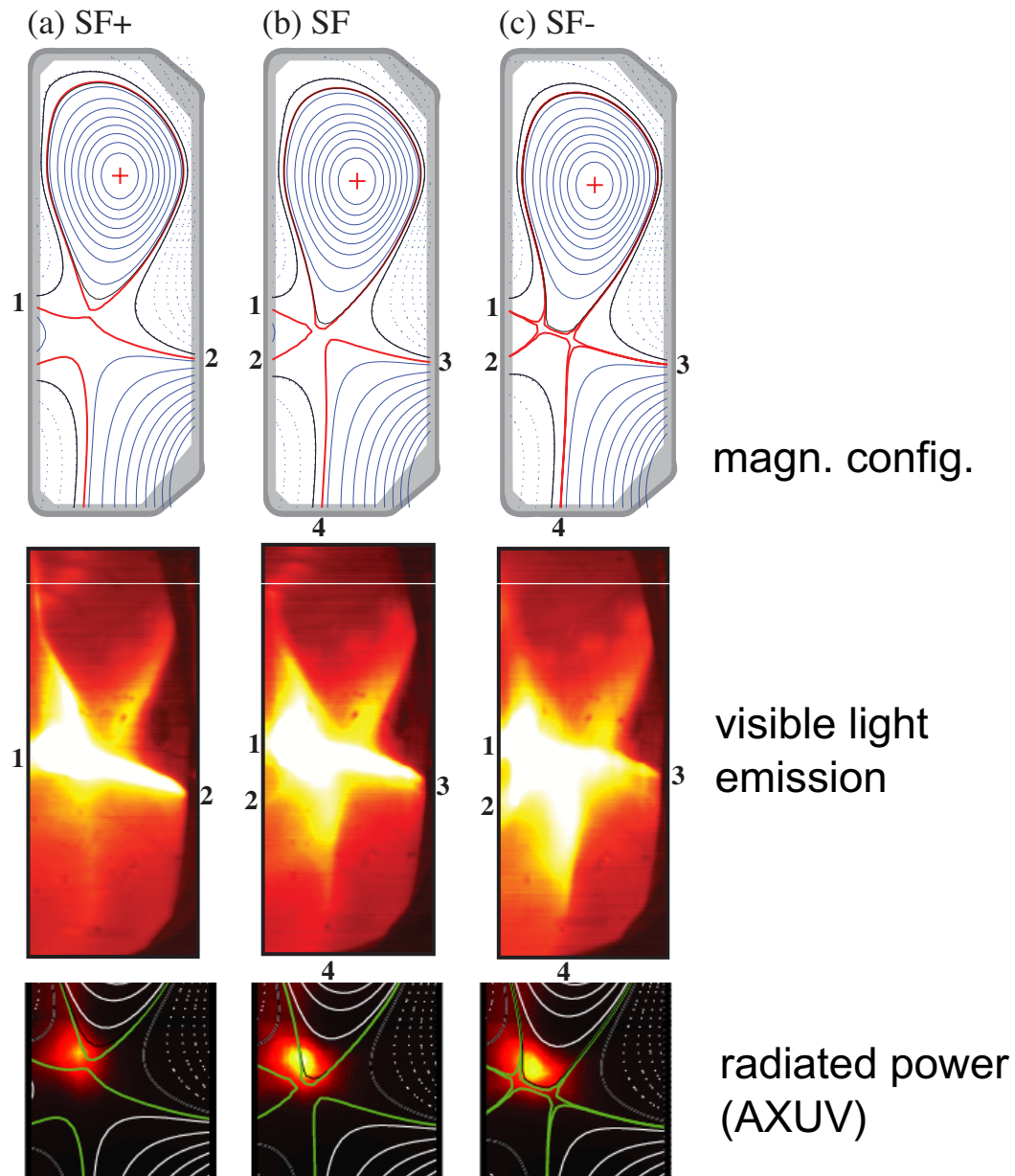
Piras PPCF 09



- **larger magn. flux expansion**
(factor 2-3 over standard X-point)
 \rightarrow heat load reduction
- **increased magnetic shear**
 \rightarrow improves edge ideal stab.,
change ELM and
H-mode properties?

Pitzschke: next talk 502 on
edge ideal-MHD stability

Snowflake divertor concept first tested in TCV

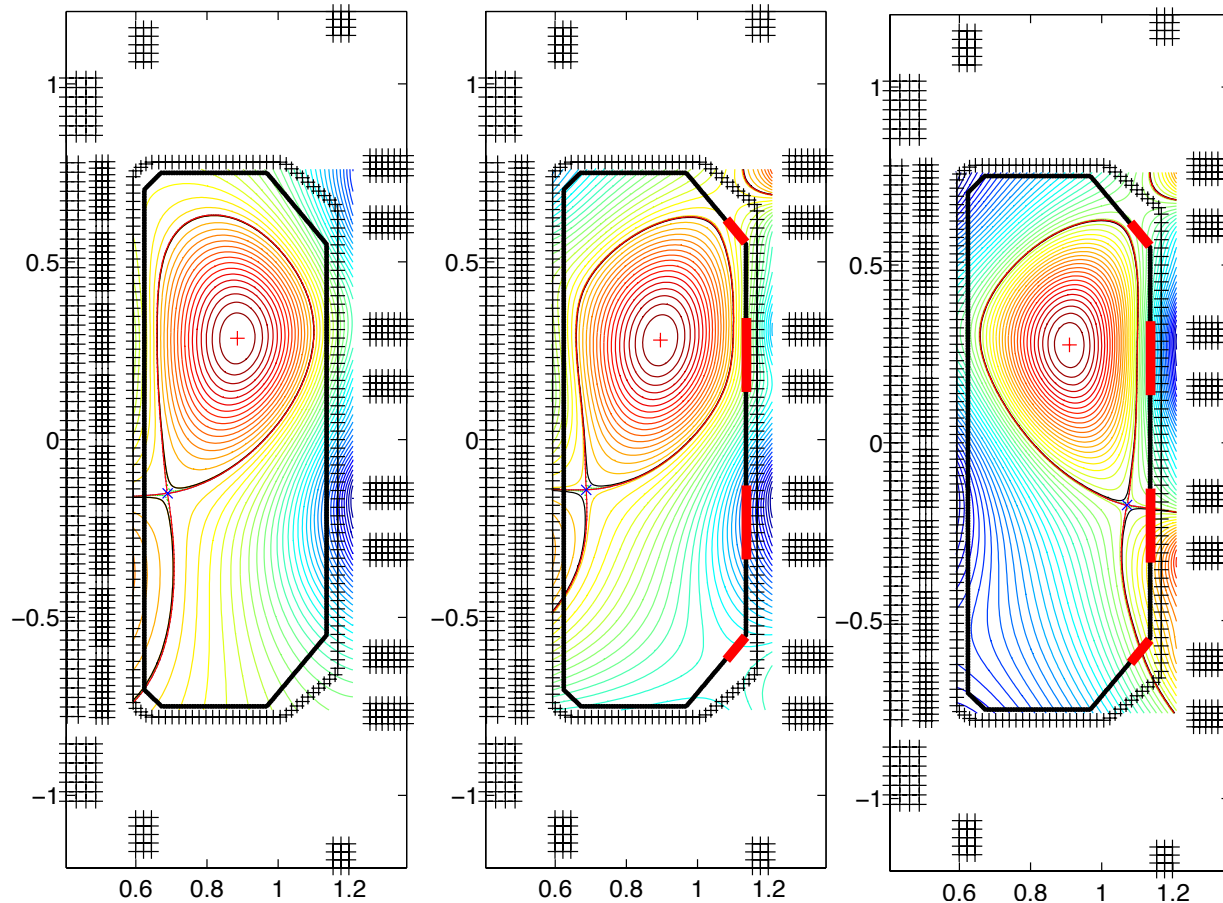


Investigation of SF and neighbour configurations (SF+, SF, SF-) :

- distrib. of power on div. legs
- distrib. of power on the walls
- maximisation of radiation power
- H-mode properties ...

Piras PPCF 09

5.2 Towards H-mode at negative triangularity



Medvedev EPS09

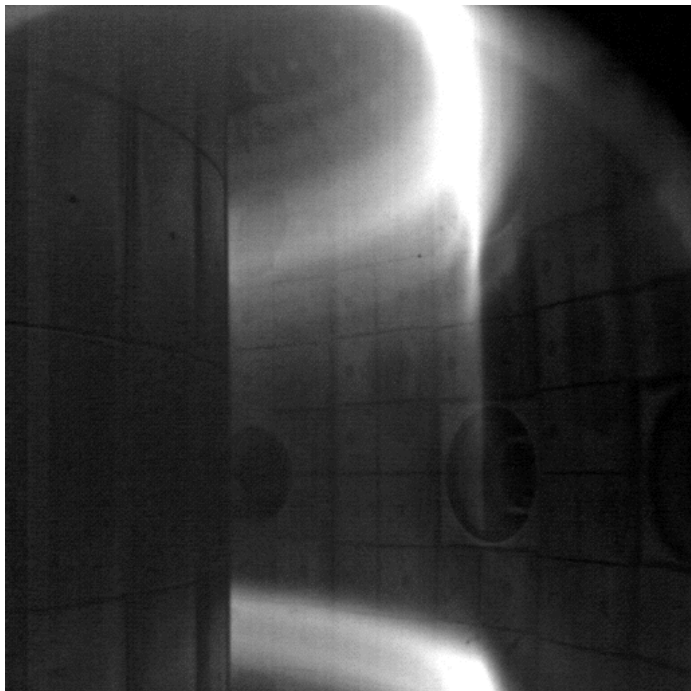
To learn more about H-mode properties by extending to negative triangularity:

from MHD modeling:
pedestal seems lower

- smaller ELM regimes?
- β -limit & RWM studies at lower power with ECH X3

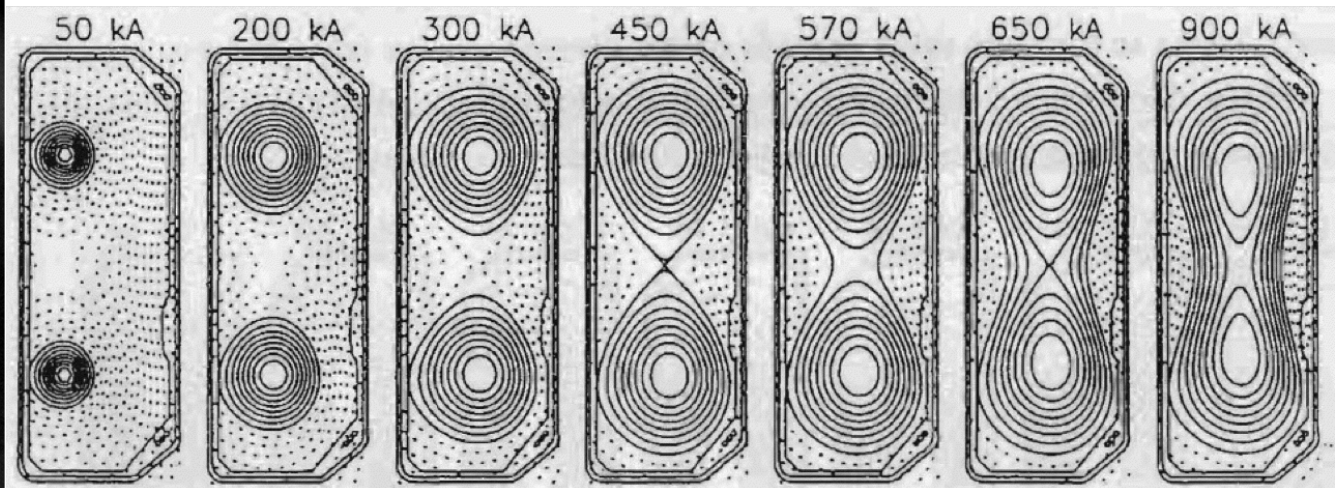
5.3 Doublet configuration: an innovative shape, a divertor concept ?

- ... on the search of configurations
- with good confinement and enhanced edge radiation properties
 - quiescent edge in H-mode scenarios...



double ECH breakdown
using 2 gyrotron beams
(Piras July 2009)

a possible scenario (model calculation)



... under development at TCV

TCV, a flexible, variable shape tokamak
to develop new shapes
to improve models

for concepts improvement
and for ITER

Thank you

CRPP-Lausanne talks / Plasma Physics Session

- **14h30, 502**

Pitzschke

MHD stability calculations for H-mode plasmas
with Snowflake divertor configuration

- **15h00, 534 (repl. 504)**

de Meijere

Measurements of electron density fluctuations in TCV

- **17h00, 508**

Furno

Turbulence and transport in TORPEX

- **17h15, 509**

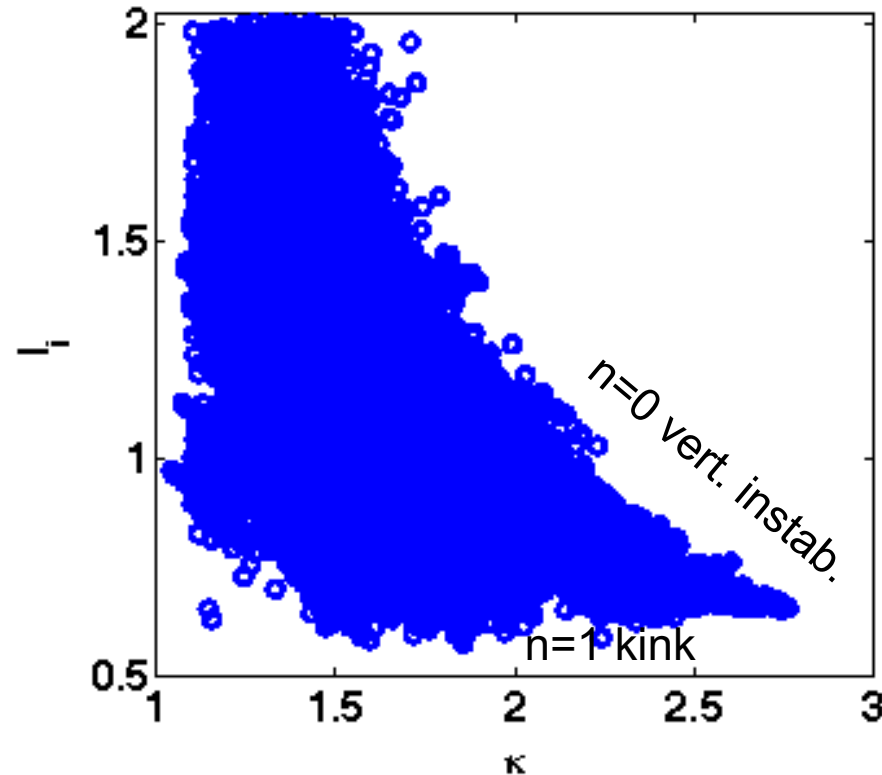
Federspiel

Observation of a critical pressure gradient
for the stabilization of interchange modes in TORPEX

SPARES

Stability limits (Ohmic)

Using high I_N , elongation up to $\kappa=2.8$



Operation at high κ limited by

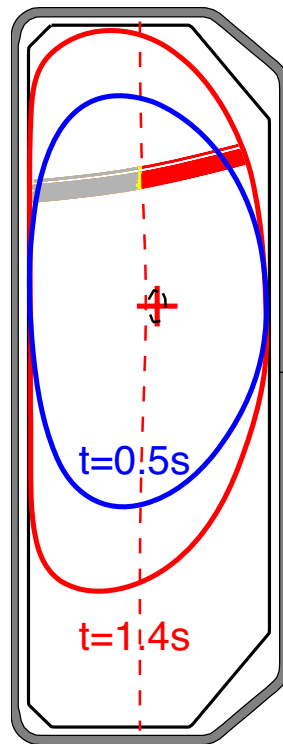
$n=0$ vertical instability \rightarrow broad $j(r)$ required: low- q , high- I_N + fast int. coils

$n=1$ external kink \rightarrow high current limit (I_N & β -limit)

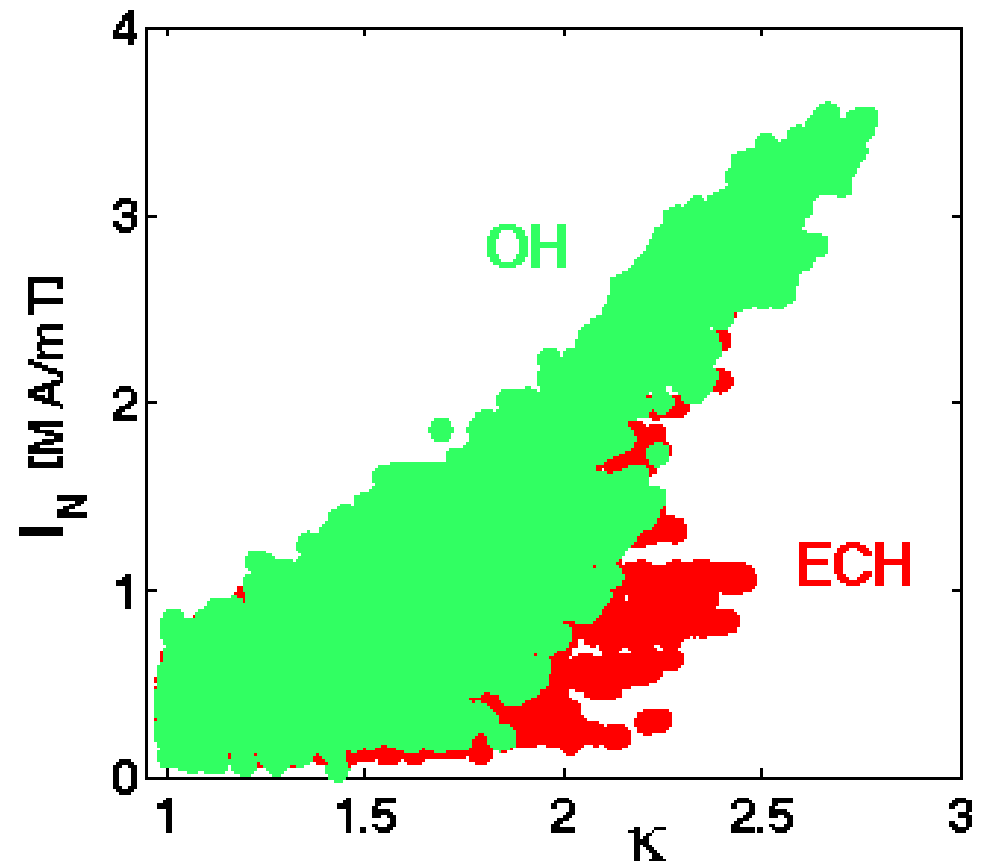
Hofmann PPCF01

High elongation stability limits (off-axis ECH)

At low I_N , vertical stability requires broadening the current profile, done using off-axis ECH, allows reaching $\kappa=2.5$ at $I_N \sim 1$

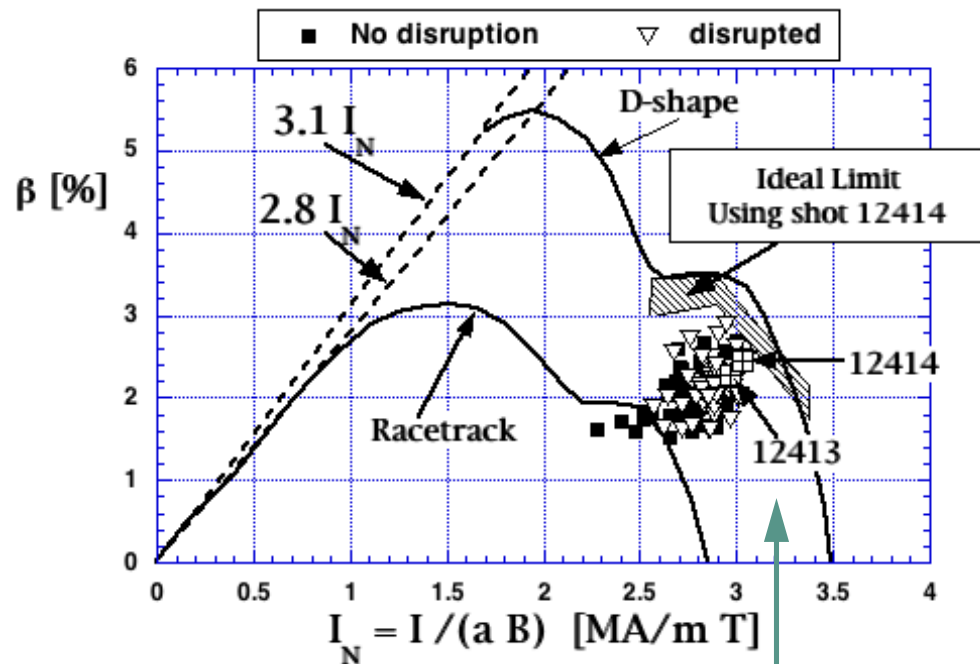


off-axis ECH power
at constant quadrupole field
—> κ increase



Pochelon NF01, Camenen NF07, Paley PPCF07

Current limit at high κ



Ideal MHD predicts correctly the current limit

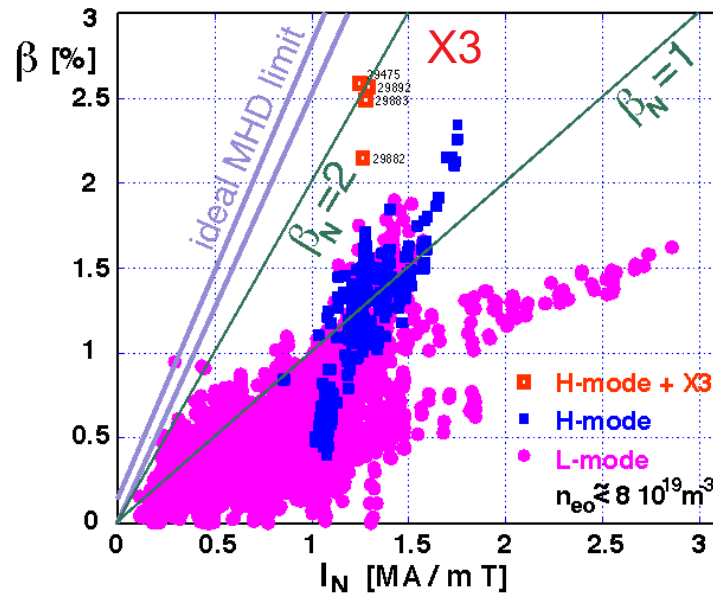
as often,
resistive modes ($m/n = 4/3, 3/2, 2/1$)
are found just below the ideal limit

current limit at high κ

shape influences the β -limit

Hofmann PRL97

β -limit at high elongation κ ...



However, more power is needed to test the β -limit in various TCV shapes
—> foreseen to double the installed X3 power

β reached at $\kappa \sim 1.6$
with 1.5 MW X3

TCV provides the highest β -values in the inter-machine spontaneous rotation database (J.Rice NF07)

Hofmann PPCF01

Porte NF07, Alberti JoPh05, Pochelon SMP05